

$$\rho = \frac{Q}{V} \quad \vec{J} = \rho \vec{u}$$



in frame of charge  $\rho_0 = \frac{Q}{V_0}$  but  $V = \frac{V_0}{\gamma}$  (by length cont.)

$$\text{so } \rho = \gamma \rho_0 \quad \vec{J} = \gamma \rho_0 \vec{u}$$

$$= \frac{\rho_0}{c} c \vec{u} = \rho_0 \vec{v}$$

recall  $c = \frac{1}{\epsilon_0 \mu_0}$

$$\rho = \rho_0 \gamma$$

$$J^\mu = \rho_0 \gamma \begin{pmatrix} c \\ J_x \\ J_y \\ J_z \end{pmatrix}$$

$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}$

Consider the continuity equation:  $\partial_\mu J^\mu = \left( \frac{1}{c} \frac{\partial}{\partial t} \frac{\partial}{\partial x} \frac{\partial}{\partial y} \frac{\partial}{\partial z} \right) \begin{pmatrix} c\rho \\ J_x \\ J_y \\ J_z \end{pmatrix}$

$$\text{note: } \vec{\nabla} \cdot (\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}) = 0$$

$$\mu_0 \vec{\nabla} \cdot \vec{J} + \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \cdot \vec{E})}{\partial t} = 0$$

$$\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$$

$$\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$$

$$= \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0 \quad \text{by continuity}$$

NOTE

next 7 pages are not  
for 3504, but included  
for completeness

Without proof:  
(shown in text)

$$\partial_\nu F^{\mu\nu} = \mu_0 J^\mu$$

$$\partial_\nu G^{\mu\nu} = 0$$

Maxwell's Eqs

$$\Rightarrow \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\Rightarrow \vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

recall:  $\vec{\nabla} \cdot \vec{B} = 0 \Rightarrow \vec{B} = \vec{\nabla} \times \vec{A}$  "vector potential"

$$\text{so } \vec{\nabla} \times \vec{E} = -\frac{\partial}{\partial t} (\vec{\nabla} \times \vec{A}) = -\vec{\nabla} \times \frac{\partial \vec{A}}{\partial t}$$

$$\text{so } \vec{\nabla} \times \left( \vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0 \Rightarrow \vec{E} + \frac{\partial \vec{A}}{\partial t} = -\vec{\nabla} \phi$$

$$\text{so } \vec{E} = -\left( \vec{\nabla} \phi + \frac{\partial \vec{A}}{\partial t} \right)$$

recall  $\vec{A}' = \vec{A} + \vec{\nabla} f$  } gauge transformations

$$\phi' = \phi - \frac{\partial f}{\partial t}$$

$$\text{ie: } \vec{B}' = \vec{\nabla} \times (\vec{A} + \vec{\nabla} f) = \vec{\nabla} \times \vec{A} + \vec{\nabla} \times \vec{\nabla} f = \vec{B}$$

$$\vec{E}' = -\left( \vec{\nabla} \left( \phi - \frac{\partial f}{\partial t} \right) + \frac{\partial}{\partial t} (\vec{A} + \vec{\nabla} f) \right)$$

$$= -\left( \vec{\nabla} \phi + \frac{\partial \vec{A}}{\partial t} \right) = \vec{E}$$

pick "Lorentz" gauge

$$\vec{\nabla} \cdot \vec{A} = -\frac{1}{c^2} \frac{\partial \phi}{\partial t}$$

so wave eqs become:

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J}$$

$$\nabla^2 \phi/c - \frac{1}{c^2} \frac{\partial^2 \phi/c}{\partial t^2} = -\mu_0 c \rho$$

(show with things like  $\vec{\nabla} \times (\vec{\nabla} \times \vec{A}) \dots$ )

but this is just

$$\underbrace{\partial_\mu}_{\text{scalar}} J^\mu A^\nu = -\mu_0 \underbrace{J^\nu}_{\text{4-vector}}$$

$$\Rightarrow A^\nu = \begin{pmatrix} \phi/c \\ A_x \\ A_y \\ A_z \end{pmatrix}$$

is a 4-vector  
(forming  $A^\nu$  can get  $\vec{E}$  &  $\vec{B}$  in diff gauges)

3. Without proof  
(done in book)

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

$$= \frac{\partial A^\nu}{\partial x^\mu} - \frac{\partial A^\mu}{\partial x^\nu}$$

↳ notice anti-symmetric

Recap:

4-vectors  
we've seen

$$x^\mu = \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$$

$$dx^\mu = \begin{pmatrix} c dt \\ dx \\ dy \\ dz \end{pmatrix}$$

$$q^\mu = \begin{pmatrix} \gamma c \\ \gamma v^1 \\ \gamma v^2 \\ \gamma v^3 \end{pmatrix}$$

$$p^\mu = \begin{pmatrix} E/c \\ p_x \\ p_y \\ p_z \end{pmatrix}$$

$$J^\mu = \begin{pmatrix} c\rho \\ J_x \\ J_y \\ J_z \end{pmatrix}$$

note  $p^\mu = m q^\mu$

$$\partial^\mu = \frac{\partial}{\partial x_\mu} = \begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial t} \\ -\frac{\partial}{\partial x} \\ -\frac{\partial}{\partial y} \\ -\frac{\partial}{\partial z} \end{pmatrix}$$

$$A^\mu = \begin{pmatrix} \phi/c \\ A_x \\ A_y \\ A_z \end{pmatrix}$$

Lorentz scalars  
we've seen  
(same in all frames)

$$x^\mu x_\mu = -c^2 t^2 + x^2 + y^2 + z^2 = -c^2 t^2 + \vec{x}^2$$

$$dx^\mu dx_\mu = -c^2 dt^2 + dx^2 + dy^2 + dz^2 = -c^2 dt^2 + d\vec{x}^2 = -c^2 d\tau^2$$

⇒  $d\tau$  is L!

$$q^\mu q_\mu = -c^2$$

$$p^\mu p_\mu = -\frac{E^2}{c^2} + \vec{p}^2 = -m^2 c^2$$

$$\Rightarrow E^2 = \vec{p}^2 c^2 + m^2 c^4$$

$$\partial^\mu \partial_\mu = -\frac{1}{c^2} \frac{\partial^2}{\partial t^2} + \nabla^2 \equiv \square^2$$

$$\partial_\nu J^\mu = 0$$

$$\partial_\mu A^\mu = 0 \Rightarrow \text{Lorentz gauge}$$

Contrav to  $\omega$ :  $x_\mu = g_{\mu\nu} x^\nu$

Forces:  $\vec{F} = \frac{d\vec{p}}{dt} \leftarrow 3 \text{ vector}$

$K^\mu = \frac{dp^\mu}{d\tau} \leftarrow 4 \text{ vector}$

EM:

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu = \begin{pmatrix} 0 & E_x/c & E_y/c & E_z/c \\ -E_x/c & 0 & B_z & -B_y \\ -E_y/c & -B_z & 0 & B_x \\ -E_z/c & B_y & -B_x & 0 \end{pmatrix}$$

"Electro-Magnetic Field Tensor"

Maxwell's Eqs:  $\partial_\nu F^{\mu\nu} = \mu_0 J^\mu$   
 $\partial_\nu G^{\mu\nu} = 0$

(Minkowski) Lorentz Force:  $K^\mu = q F^{\mu\nu} v_\nu$

Wave Eqn:  $\partial_\mu \partial^\mu A^\nu = -\mu_0 J^\nu$  (in Lorentz gauge  $\partial_\mu A^\mu = 0$ )  
 $\square^2 A^\nu = -\mu_0 J^\nu$

Switching frames

$$\bar{X}^\mu = \Lambda^\mu_\nu X^\nu \quad X^\nu \text{ an arbitrary 4-vector}$$

$$\bar{X}_\mu = \Lambda_\mu^\nu X_\nu$$

ie:  $\bar{K}^\mu = \Lambda^\mu_\nu K^\nu$   
 $\bar{v}^\mu = \Lambda^\mu_\nu v^\nu$

Special cases  
ie dt instead  
of t dt

normal forces =  $\frac{d\vec{p}}{dt}$   
ie boost in x

$$\bar{F}_y = \frac{F_y}{\gamma(1 - \beta u_x/c)}$$

$$\bar{F}_z = \frac{F_z}{\gamma(1 - \beta u_x/c)}$$

$$\bar{F}_x = \frac{F_x - \beta(\vec{v} \cdot \vec{F})/c}{(1 - \beta u_x/c)}$$

normal velocities =  $\frac{dx}{dt}$   
ie boost in x

$$\bar{u}_x = \frac{u_x - v}{(1 - v u_x/c^2)}$$

$$\bar{u}_y = \frac{u_y}{\gamma(1 - v u_x/c^2)}$$

$$\bar{u}_z = \frac{u_z}{\gamma(1 - v u_x/c^2)}$$

$$F^{\alpha\beta} = \Lambda^{\alpha}_{\mu} \Lambda^{\beta}_{\nu} F^{\mu\nu} = \Lambda^{\alpha}_{\mu} F^{\mu\nu} (\Lambda^{\dagger})^{\beta}_{\nu} = \underbrace{\Lambda F \Lambda^{\dagger}}_{\text{as matrices}}$$

ie boost in x

$$\bar{E}_x = E_x$$

$$\bar{E}_y = \gamma(E_y - vB_z)$$

$$\bar{E}_z = \gamma(E_z + vB_y)$$

$$\bar{B}_x = B_x$$

$$\bar{B}_y = \gamma(B_y + \frac{vE_z}{c^2})$$

$$\bar{B}_z = \gamma(B_z - \frac{vE_y}{c^2})$$

$$\bar{F} = \left( \Lambda \right) \left( F \right) \left( \Lambda^{\dagger} \right)$$

$\Lambda$  options: rotate about x:

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & c & -s \\ 0 & 0 & s & c \end{pmatrix}$$

$$s \equiv \sin \theta$$

$$c \equiv \cos \theta$$

$$y: \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & c & 0 & -s \\ 0 & 0 & 1 & 0 \\ 0 & s & 0 & c \end{pmatrix}$$

orthogonality  $\Lambda^{\alpha}_{\beta} \Lambda^{\beta}_{\alpha} = \delta^{\alpha}_{\alpha}$

$$z: \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & c & -s & 0 \\ 0 & s & c & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

boost in x:

$$\begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

note: under boosts components of  $\vec{E} + \vec{B}$  get mixed up

$$y: \begin{pmatrix} \gamma & 0 & -\beta\gamma & 0 \\ 0 & 1 & 0 & 0 \\ -\beta\gamma & 0 & \gamma & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$z: \begin{pmatrix} \gamma & 0 & 0 & -\beta\gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\beta\gamma & 0 & 0 & \gamma \end{pmatrix}$$

In fact, it probably confuses more than enlightens since in a matrix it is hard to tell if the indices are supposed to be superscript or subscript. In fact for anything above rank 2 tensor, it's not possible to write in "matrix form".

Correction to H.O. : 1) in Mathematics it should read  
 transform = lambda . f. Transpose [lambda]  
 (makes no difference for booleans, but does affect relations)  
 2) on last page.

$$\vec{F} = (\Delta)(F)(\Lambda^t)$$

note:  $\partial_\mu J^\mu = \left( \frac{1}{c} \frac{\partial}{\partial t} \quad \frac{\partial}{\partial x} \quad \frac{\partial}{\partial y} \quad \frac{\partial}{\partial z} \right) \begin{pmatrix} c\rho \\ J_x \\ J_y \\ J_z \end{pmatrix} = \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0$  (by continuity)

As Acron noted in class. note: + since  $\partial_\mu = \frac{\partial}{\partial x^\mu}$

consider  $\eta^\mu \eta_\mu = (c\gamma, n_1, n_2, n_3) \begin{pmatrix} -c\gamma \\ n_1 \\ n_2 \\ n_3 \end{pmatrix} = -c^2 \gamma^2 + \vec{n}^2 = -c^2$  in frame of particle

we use  $g = \begin{pmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$  but more common to use  $\begin{pmatrix} 1 & & & \\ & -1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$

$\Rightarrow \eta^\mu \eta_\mu = (c\gamma, n_1, n_2, n_3) \begin{pmatrix} c\gamma \\ -n_1 \\ -n_2 \\ -n_3 \end{pmatrix} = c^2 \gamma^2 - \vec{n}^2 = c^2$  in frame of particle

does this make sense? In frame of particle, it moves forward in TIME ( $\tau$ ) at  $c$ . In another frame as  $\vec{n} \rightarrow$  large  $\gamma \rightarrow$  large such that  $c^2 \gamma^2 - \vec{n}^2 = \text{constant} = c^2 (LI)$

A particle moving past you w/ velocity  $\vec{u}$  also has its clock running slow, but  $dt = \gamma d\tau$  so for you, an increase in the spatial part of the 4-velocity  $\vec{v} = \frac{d\vec{x}}{dt}$  also gives an increase in the time component  $v^0 = \frac{dt}{d\tau} = \gamma$

!! ADVANCED !!

where else have you seen "similarity" forms?

NOT 3504

recall:  $\vec{I} \vec{\omega}_k = \lambda_k \vec{\omega}_k \Rightarrow \vec{\omega}_k, \lambda_k$   
 (from lin. alg. or mechanics)

form  $\vec{A} = \begin{pmatrix} \omega_1 & \dots \\ \omega_2 & \dots \\ \omega_3 & \dots \end{pmatrix}$  note: normalized eigenvectors

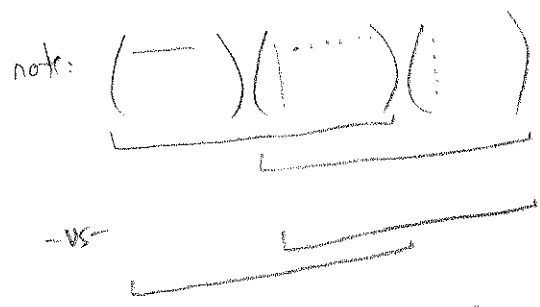
then  $\vec{A}^{-1} = \vec{A}^t$

so  $\vec{A} \vec{I} \vec{A}^t \vec{A} \vec{\omega}_k = \lambda \vec{A} \vec{\omega}_k$   
 $\vec{I}' \vec{\omega}' = \lambda \vec{\omega}'$

$\vec{I}' = \vec{A} \vec{I} \vec{A}^t$   
 $\vec{\omega}' = \vec{A} \vec{\omega}$

$\vec{x} = \vec{A}^{-1} \vec{v} (= \vec{A}^t \vec{v})$   
 check:  $\vec{x}_1 = \begin{pmatrix} \omega_1 & \omega_2 & \omega_3 \\ \vdots & \vdots & \vdots \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$

so we have  $\vec{I}' = \vec{A} \vec{I} \vec{A}^t$  (since  $\vec{A}$  is orthonormal)  $= \begin{pmatrix} \omega_{1x} & & \\ & \omega_{2y} & \\ & & \omega_{3z} \end{pmatrix}$



ie: order in which multiplication is carried out doesn't matter  
 ie  $(AB)C = A(BC)$

do this so that rows & columns are next to each other + can write like matrix mult.

$F^{\alpha\beta} = \Delta^{\alpha}_{\mu} F^{\mu\nu} \Delta^{\beta}_{\nu}$   
 $= \Delta^{\alpha}_{\mu} F^{\mu\nu} \Delta^{\beta}_{\nu} = \Delta^{\alpha}_{\mu} F^{\mu\nu} \Delta^{\beta}_{\nu} = (\Delta)(F)(\Delta^t)$

ie: boost in x:  $F = \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} F \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$

rotate about x:  $F = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & c & -s \\ 0 & 0 & s & c \end{pmatrix} F \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & c & s \\ 0 & 0 & -s & c \end{pmatrix}$

so...

$K^{\mu\nu} = g F^{\mu\lambda} \Omega_{\lambda}^{\nu}$

of equiv.  $K^{\mu\nu} = g F^{\mu\lambda} \Omega_{\lambda}^{\nu}$

$\begin{pmatrix} \dots \\ \dots \\ \dots \end{pmatrix} \begin{pmatrix} \dots \\ \dots \\ \dots \end{pmatrix}$

!! ADVANCED !! what about Lagrangian?

$L = \frac{-mc^2}{\gamma} + q\vec{v} \cdot \vec{A} - q\phi$  will do (some "guess")

X Comp

$\frac{\partial L}{\partial v_x} = \gamma m v_x + q A_x = \vec{p}_x$  "canonical momentum"

$\frac{d}{dt} \frac{\partial L}{\partial v_x} = \frac{d}{dt} (\gamma m v_x) + q \left( \frac{\partial A_x}{\partial t} + \frac{\partial A_x}{\partial x} v_x + \frac{\partial A_x}{\partial y} v_y + \frac{\partial A_x}{\partial z} v_z \right)$

$\frac{\partial L}{\partial x} = -q \frac{\partial \phi}{\partial x} + q \left( v_x \frac{\partial A_x}{\partial x} + v_y \frac{\partial A_y}{\partial x} + v_z \frac{\partial A_z}{\partial x} \right)$  but:  $\frac{d}{dt} \frac{\partial L}{\partial v_x} = \frac{\partial L}{\partial x}$

so:  $\frac{d}{dt} (\gamma m v_x) = -q \left( \frac{\partial \phi}{\partial x} + \frac{\partial A_x}{\partial t} \right) + q \left( v_y \left[ \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right] + v_z \left[ \frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} \right] \right)$   
 $= q E_x + q (\vec{v} \times \vec{B})_x = \frac{d p_x}{dt} = F_x$  ✓

$H = \vec{p} \cdot \vec{v} - L$

$= \frac{\vec{p} \cdot \vec{v} c}{\sqrt{m^2 c^2 + d^2}} + \frac{m^2 c^3}{\sqrt{m^2 c^2 + d^2}} - \frac{q \vec{A} \cdot \vec{v} c}{\sqrt{m^2 c^2 + d^2}} + q \phi$

$\gamma m \vec{v} = \vec{p} - q \vec{A}$

$\frac{m^2 v^2}{1 - \frac{v^2}{c^2}} = d^2$

$= \frac{[(\vec{p} - q\vec{A})^2 + m^2 c^2] c}{\sqrt{(\vec{p} - q\vec{A})^2 + m^2 c^2}} + q \phi$

$m^2 v^2 = d^2 - \frac{d^2 v^2}{c^2}$

$= \sqrt{(\vec{p} - q\vec{A})^2 c^2 + m^2 c^4} + q \phi$

$v^2 (m^2 + \frac{d^2}{c^2}) = d^2$

recall:  $E = \sqrt{p^2 c^2 + m^2 c^4}$   
 so for EM  $\vec{p} \rightarrow \vec{p} - q\vec{A} + \text{odd } q\phi$

$\vec{v} = \frac{\vec{v} c}{\sqrt{m^2 c^2 + d^2}}$

$\frac{v^2}{c^2} = \frac{d^2}{m^2 c^2 + d^2}$

$1 - \frac{v^2}{c^2} = \frac{m^2 c^2 + d^2 - d^2}{m^2 c^2 + d^2}$

$\gamma = \frac{\sqrt{m^2 c^2 + d^2}}{m c}$